# Edge state and unconventional magnetism of nanographene/nanographite

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nanographite  $\pi$ -electron system with open edges



fullerenes carbon nanotubes closed π-electron system



nanographite (nanographene)

localized  $\pi$ -spins

enhanced magnetism





non-Kékule str. nonbonding  $\pi$ -state (s=172) Yamabe et al. Fujita et al.



zigzag edge

armchair edge











ferromagnetic

edge state nonbonding π-state

large density of states (DOS) spin paramagnetism

contrast to bulk graphite

#### extra $\pi$ -bond connected to the carbon atom at a zigzag edge

edge states with localized spins appear at -2/3 < k < 2/3

Klein edge

K. Kusakabe, et al. 2001

#### combination of Fujita's and Klein's edges



completely localized edge state appears at  $E_{\rm F}$ 

K. Kusakabe, et al. 2001

nano-magnetism

#### Kusakabe et al.

#### Carbon-only ferromagnet





spatially maldistributed ferromagnetism

#### # Nanographene

# single nanographene sheet
 nanodiamond
 electrophoretic technique + heat-treatment
# nanographene ribbon
resonance Raman experiments
# electronic structure of edge state
 at edges hydrogen-terminated

*#* Controllable nanoscopic magnetism

gas-adsorption induced magnetic switch

# single nanographene sheet

#### SEM and AFM images of diamond nanoparticles.





nanodiamond particles (ca. 5 nm) deposited by electrophoretic technique spherical shape with particle sizes larger than those observed by TEM absorbed solvent molecules on the surface of particles

#### nanographene and STM analysis



STM images after heat-treatment at 1600 °C

nanographene > flat single layer sheet mean in-plane size of 10nm

#### nanographene on HOPG substrate and STM Analysis





#### cross-sectional profile

# nanographene ribbon observed by resonance Raman experiments

#### AFM image of single nanographene ribbon

#### single sheet of nanographene ribbon at a step edge



#### Resonance Raman experiments with polarized light



small nanographene ribbon can be easily heated by laser beam

edge states observed by STM/STS and analyzed by tight-binding calculations

#### zigzag edge electronic state of graphene edges 150 (NYU) experimental evidence of edge edge state state dl/dVs 31 Zigzag 0 0.5 0 -0.5 Armchair Vs (V) armchair edge Armchair 200dl/dV (nA/V) π\* Armchair 0 -0.5 0.5 0 Vs (V)

zigzag edge: short and defective, energetically unstable

(a)

3 zigzag carbon sites

finite length zigzag edge



(b)

localized edge state

tight binding calculation

edge state

#### electron confinement effect in zigzag edges

#### edge-state-absent site at zigzag edge (small local density of states (LDOS))



# magnetic switch of nanographite

#### Activated carbon fibers (ACFs)



#### 3D random network of nanographite domains



edge-state spins

adsorption of guest molecules

#### Gas physisorption $(O_2, N_2, H_2O \text{ etc.})$ in ACF nanopore space



guest molecules -> effective pressure inter-nanographene layer distance reduced What effect on magnetism ? nanographite + H<sub>2</sub>O

#### Comparison between H<sub>2</sub>O adsorption isotherm and magnetic susceptibility



# physisorption isotherm
 adsorption threshold
 at P/P<sub>0</sub> ~ 0.5
 (hydrophobic pore)

# magnetic susceptibility
 drops at P/P\_0 > 0.5

ON/OFF magnetic switch

#### effective pressure induced by H<sub>2</sub>O guests



T.Suzuki et al., Carbon(1988)

change in magnetism (high spin - low spin transition)



inter-graphene interaction  $J \sim t^2/U$  enhanced effective magnetic moment reduced

controllable nanoscopic magnetism in nanographite

#### Analysis by the Hubbard type model

intralayer transfer int.: t
on-site Coulomb int.: U

# closed shell electron case:

increase of interactions→ increase of total magnetic moment

# open shell electron case:

increase of interactions→ decrease of total magnetic moment

→ in agreement with experiments



K. Harigaya, J. Phys.: Condens. Matter <u>13</u>, 1295 (2001);
K. Harigaya, Chem. Phys. Lett. <u>340</u>, 123 (2001);
K. Harigaya and T. Enoki, Chem. Phys. Lett. <u>351</u>, 128 (2002).

#### Ar adsorption in micropores



magnetization affected well above the boiling point

Ar guest atoms (diamagnetic) adsorbed well above the boiling point

#### zero-field muon spin relaxation (ZF- $\mu$ SR) in Ar



due to the dipolar field



## Summary

Nanographene/Nanographite

importance of geometry of edges

non-bonding  $\pi$ -electron state (edge state)

STM/STS observations for well defined edges

nanoscopic magnetism

spin glass magnetic switch gas sensor

# Nanographene-based molecular devices

zigzag edge armchair edge



# electron beam Ithography magnetic line

#### chemical modifications

CH itinerant magnetism
CH<sub>2</sub> localized magnetism
CF nonmagnetic
C=O conducting line

future promising molecular devices (compared with nanotubes)

#### Contributors

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Nanographene ribbon



#### magnetic nanocarbons with Klein and Fujita edges



#### nanocarbon-based ferromagnetism

F(or H)-terminated edge graphite-diamond interface

F

K. Kusakabe, et al. 2001

#### triangular lattice image on a particle



★same A-B stacking mode with HOPG substrate
★small irregularity

#### **Electronic properties**

#### •graphene sheet-substrate interaction

interlayer distance 0.35-0.37 nm



interlayer resonance integral  $\gamma_1 = 0.20 - 0.29 eV$ (bulk  $\gamma_1 = 0.39 eV$ )

#### +quantum size effect

 $N \sim 3000$  largest polycyclic aromatic molecule band width  $W = 6\gamma_0$ in-plane resonance integral  $\gamma_0 = 3.16$  eV energy discreteness  $N/W \sim 75$  K

+edge state effect

scanning tunneling spectroscopy experiment



dispersed LDOS near the defect point at which an armchair line is added

#### two zigzag edge atoms added to armchair edge



 $(4.3 \times 4.3 \text{ nm}^2)$ 

A

increase by two armchair lines

(4.5×4.5 nm<sup>2</sup>)

B

arrays of bright spots near defects run in different directions



calculated LDOS near the defect point



the difference comes from one extra-carbon atom

#### effect of inter-graphene layer interaction on the edge state

#### 2D LDOS mapping



#### effect of various physisorbed molecules on nanomagntism



# magnetic switch effect created at the solid-liquid transition of $Br_2$ guests



magnetoresitance with paramgnetic  $O_2$  molecules



#### $O_2$ molecules S=1



wave function overlap with graphitic  $\pi$ -electron strong exchange int.

chemical bond with S=0 singlet state

physisorbed weak dipole-dipole int.

exchange int. between  $O_2$  spins and edge-state spins

internal field from O<sub>2</sub> molecules ~20 K

#### Zero field muon spin relaxation ( $\mu$ SR) in Ar atmosphere



Ar pressure: ~ 1 atm at 87 K (B.P.)

•: 84 K ⊽: 95 K △: 294 K

Freezing P. 84 K Boiling p. 87 K

increase of the static component is prominent below the F.P. of Ar

freezing of Ar condensed in the nanopores modification of the magnetism of the edge-state spins

# nano-graphite network

☆Coulomb blockade effect
 ☆spin-glass
 Activated carbon fibers
 3D nano-graphite network

heat treatment

insulator-metal transition

#### Heat-treatment effect on conductivity



semiconductive region conductivity of ACFs and iodine-doped ACFs

Coulomb gap variable-range hopping conduction

e<sup>2</sup> Er

# charging effect





$$\sigma = \sigma_0 \exp\left[ (T_0 / T)^{1/2} \right]$$
$$T_0 = \frac{6e^2}{\pi k_B} \frac{1}{4\pi\varepsilon} \frac{1}{\alpha^{-1}}$$
$$\varepsilon: \text{ dielectric constant}$$
$$\alpha^{-1}: \text{ localization length}$$

#### Heat-treatment effect on magnetic susceptibility



HTT <  $P_c$ Curie-Weiss behavior with localized spins and negative Weiss temperature

HTT >  $P_c$ less temp. dependent enhanced diamagnetism

### Magnetism around the percolation threshold region

![](_page_49_Figure_1.jpeg)

antiferromagnetic ordering? negative Weiss temperature

#### Field cooling effect

large field cooling effect around the MI threshold

![](_page_50_Figure_2.jpeg)

exchange interaction  $|\sqrt{\langle \Delta J^2 \rangle}/\langle J \rangle| \sim 0.8$ random distribution

#### ACFs

![](_page_51_Figure_1.jpeg)

# 3D random network of nanographite domains
# nanopores
# edge-state spins

# Huge helium condensation in micropores

#### Activated carbon fibers (ACF) 3D random network of nanographites

![](_page_53_Figure_1.jpeg)

![](_page_53_Picture_2.jpeg)

nanographite; metallic domain

3-4 graphene sheets with in-plane size 2-3nm

 localized spins of edge origin → several spins / nanographite
 micropores ~ 1-2nm, networked

(Helium in micropores) ACF; specific surface areas  $\sim 3000 \text{ m}^2/\text{g}$ 🛪 gas adsorption He, Ne, Ar, H2, N2, O2 \* ESR spin-lattice relaxation saturation technique edge-localized spins probe for guest-host int. ħω spin system  $T_1$ lattice

![](_page_55_Figure_0.jpeg)

effective pressure in micropores at room temp.

large condensation of guest gaseous species

![](_page_55_Figure_3.jpeg)

extremely large condensation ~ 10<sup>4</sup> times

He atom, exceptional nature

#### ESR saturation curves with guest gases

![](_page_56_Figure_1.jpeg)

independent of internal degree of freedom rotation / vibration

spin-lattice relaxation rate  $T_1^{-1}$ is accelerated by gas uptake remarkable enhancement in He atmosphere spin-lattice relaxation rate  $T_1^{-1}$ 

![](_page_57_Figure_1.jpeg)

![](_page_57_Figure_2.jpeg)

1/T<sub>1</sub>=nσυ
n: density of helium atoms
σ: cross section related to the spin-flop process
υ: velocity of an He atom

![](_page_57_Figure_4.jpeg)

$$\sqrt{\langle v^2 \rangle} = \sqrt{\frac{3RT}{M}}$$
$$1/T_1 \sim \sqrt{T}$$

 $\frac{1}{T_1} = \frac{6.05}{9(2\pi)^{1/2}h^4} n M^{3/2} \frac{p}{\Delta E} \frac{(e^2 IJ)^4}{R_0^6} w(k_BT)^{1/2}$  *M*: mass of He atom, *R*o: min.dist.(He atom and edge spin)  $\lambda$ : spin-orbit int. of carbon p state, D: difference (gr. and ex. states of C)  $\Delta E$ : difference (<sup>1</sup>P and <sup>1</sup>S states of He atom)

 $p = \lambda / \Delta$   $I = \int \Psi_{1s} \Psi_{2pz} dr$   $J = \int \Psi_{2s} \Psi_{2pz} dr$   $\Psi$ : He wave function

helium atoms in activated carbon fibers  $\odot$  extremely large condensation at room temperature  $\odot$  remarkable acceleration of  $T_1^{-1}$ 

![](_page_58_Figure_3.jpeg)

He atom in ultra-micropore

enhanced interaction He-nanographite

![](_page_59_Figure_0.jpeg)

nano-graphite  $\pi$ -electron system with open edges

![](_page_59_Picture_2.jpeg)

fullerenes carbon nanotubes closed π-electron system

![](_page_59_Picture_4.jpeg)

Nanographene ribbon

![](_page_60_Picture_1.jpeg)

#### combination of Fujita's and Klein's edges

![](_page_61_Figure_1.jpeg)

completely localized edge state appears at  $E_{\rm F}$ 

K. Kusakabe, et al. 2001

nano-magnetism

#### Klein edge π-bond connected to the carbon atom at a zigzag edge

![](_page_62_Figure_1.jpeg)

![](_page_62_Figure_2.jpeg)

#### edge states with localized spins appear at -2/3 ≦k< 2/3

K. Kusakabe, et al. 2001

#### magnetic nanocarbons with Klein and Fujita edges

![](_page_63_Picture_1.jpeg)

#### nanocarbon-based ferromagnetism

F(or H)-terminated edge graphite-diamond interface

K. Kusakabe, et al. 2001

# Electron wave diffractions by STM

# wave function diffraction patterns observed by STM theoretical characterization

# Variations in diffraction period

![](_page_65_Picture_1.jpeg)

![](_page_66_Figure_0.jpeg)

# tilted with respect to the substrate

potential gradient on the graphene plane

free electron model

$$\left[-\frac{\hbar^2}{2m}\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right) + V_{\text{well}}(x) - F \cdot y\right]\psi(x, y) = E\psi(x, y)$$

linear potential

# Electron density by the k p model

Hamiltonian around K-point

 $H = \begin{pmatrix} -F \cdot y & -i\gamma \frac{\partial}{\partial x} - \gamma \frac{\partial}{\partial y} \\ -i\gamma \frac{\partial}{\partial x} + \gamma \frac{\partial}{\partial y} & -F \cdot y \end{pmatrix}$ 

Solution with the infinite well

$$\Psi = 2A \left( \begin{array}{c} \sin\left(\frac{E_n x}{\gamma}\right) \sin\left[\frac{1}{\gamma} \left(\tilde{E}y + \frac{1}{2}Fy^2\right)\right] \\ -i\cos\left(\frac{E_n x}{\gamma}\right) \cos\left[\frac{1}{\gamma} \left(\tilde{E}y + \frac{1}{2}Fy^2\right)\right] \end{array} \right)$$

linear potential 13 20.5

X

electron density at the A-sublattice  $|\psi_{A}(\mathbf{R}_{A})|^{2} = 4A^{2}\sin^{2}\left(\frac{n\pi x}{d}\right)$  $\times \left\{1 + \cos[(\mathbf{K} - \mathbf{K'}) \cdot \mathbf{R}_{A} - \eta]\sin\left[\frac{2}{\gamma}\left(\tilde{E}y + \frac{1}{2}Fy^{2}\right)\right]\right\}$ 

■ 1.5-2
□ 1−1.5
■ 0.5-1
■ 0-0.5

0.5

envelope function as a long range component

Local density of states

 $\sin^2\left(\frac{n\pi x}{d}\right)\left[\operatorname{const.}+\sin\left(\frac{Fy^2}{\gamma}-\frac{2n\pi}{d}y\right)\right]$ 

-0.5

#### Contributors

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#### He pressure dependence of spin-lattice relaxation rate at room temperature

![](_page_69_Figure_1.jpeg)

enhancement of  $T_1^{-1}$ collisional process of He atom

# Present talk

#### # Nano-graphite

heat treatment of nano-diamond at 1600°C
 electrophoretic technique
 single nano-graphene sheet
 relectronic structure

# Nano-graphite network activated carbon fibers (ACF) ☆ Coulomb blockade effect ☆ spin-glass

#### # Controllable nanoscopic magnetism

water-adsorption induced magnetic switch